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Thermal Buckling and Thermal Induced free Vibration Analysis of Perforated Composite Plates: a Mathematical Model

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KEYWORDS

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ABSTRACT

This article is concerned with thermal buckling and thermal induced free vibration analyses of PCPs (perforated composite plates) with simply supported edges applying a mathematical model. The stiffness and density of PCP are defined locally using Heaviside distribution functions. The governing equations are derived based on CLPT. The present solution gives reasonable results in comparison with the few literatures. In order to inspect the structural behaviour of PCPs subjected to initial thermal loads, many parametric studies have been carried out. Results indicated that the presence of perforations has a significant effect on thermal buckling and thermal induced fundamental frequency.

1. Introduction

The increased application of light weight engineering structures has stimulated interest in the improvement of methods for the analysis of PCPs under different loadings. In order to design PCPs in thermal environments, considerations must be applied to the thermally induced vibration behavior and the possibility of thermal buckling. As the first attempt on thermal buckling of flat or initially imperfect isotropic plates, Gossard et al. [1] performed an approximate solution based on large deflection plate theory. Tauchert and Huang [2] employed the Rayleigh-Ritz method in order to solve the thermal buckling equations for symmetric angle-ply laminated plates. Thangaratnam applied [3] the linear theory and the finite element method (FEM) to solve the thermal buckling problem of composite laminated plates. It was concluded that fiber orientation, number of layers, aspect ratio, and edge conditions could influence the critical buckling temperature and mode shapes significantly. Sunt and Hsu [4] investigated the thermal buckling problem of symmetric cross-ply laminated plates applying Kirchhoff deformation theory with the inclusion of transverse shear deformation in the displacement field. The critical temperature results were acquired by the Navier solution procedure. They

revealed that the discrepancy of results obtained by the formulation with or without shear deformation components is noticeable for length to thickness ratios smaller than 20. Thronton [5] demonstrated research to review the temperature distributions in thin walled structures and thermal buckling analysis methods of isotropic and composite shells and plates. Whitney [6] examined expansional strain effects in laminated plates and derived formulations for bending and thermal buckling problems. Kant and Babu [7] applied two shear deformable finite element models based on first order shear deformation theory (FSDT) and higher order shear deformation theory (HSDT) for thermal buckling of skew fiber reinforced composite and sandwich plates. They indicated that the critical temperature values increase with an increase in skew angle and it is more pronounced in thin laminates than in thick ones. Yapici [8] investigated the thermal buckling behavior of hybrid-composite angle-ply laminated plates with an inclined crack using FEM. Pradeep and Ganesan [9] explored thermal buckling of multi-layer rectangular viscoelastic clamped sandwich plates using FEM. Duran et al. [10] acquired thermal critical buckling temperatures of composite plates with spatial varying fiber orientations using classical lamination theory (CLPT) and FEM. Li et al. [11] applied CLPT in

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order to investigate buckling and vibro-acoustic responses of the clamped composite plates in thermal environments excited by a concentrated harmonic force. Jin et al. [12] performed digital image correlation (DIC) technique to inspect thermal buckling measurement of a circular laminated composite plate under uniform temperature distribution. Through making comparisons among experimental results and those obtained by theory and nonlinear FEM analysis, they exhibited that DIC is promising for studying thermal buckling of composite structures in diverse fields. Bouazza et al. [13] studied thermal buckling of laminated cross-ply plates applying a refined hyperbolic shear deformation theory. Cetkovic [14] proposed a layerwise displacement model to analyze thermal buckling problem of composite plates and showed that the critical buckling temperature increases by increasing the aspect ratio of plate. Kalita and Haldar [15] reported the free vibration analysis of rectangular isotropic plates using a nine node isoparametric plate element in conjunction with first-order shear deformation theory contemplating the effect of rotary inertia. Perforated structures primarily have been modeled applying equivalent stiffness. A few researches [16-18] have been published to analyze structures including discontinuities using exact modeling. Takabatake [16] examined the static analysis of isotropic plates with voids using the unit step function to define the structure stiffness. Takabatake used the Galerkin method to solve the differential equation of motion. Li and Cheng [17] analyzed grid stiffened composite sandwich panels with simply supported edges subjected to lateral uniform pressure. For an orthogrid stiffened plate, they considered two material regions, the cell, and the surrounding ribs. Based on this concept, they modeled the grid shape in terms of Heaviside functions, which results in the local definition for ABD matrices. The governing equations are solved by deliberating only one component of displacement w , so the solution is limited to symmetric sandwich lay-ups. Wilson et al. [18] have operated research on elastic stability of stepped and stiffened plates. They modeled structures with variation in thickness such as single step, double stepped and latitudinal stiffened plates applying piecewise functions for thickness. The current study presents an analytical solution to thermal buckling and thermal induced free vibration analysis of PCPs taking into account uniform temperature rise and simply supported (SSSS) boundary conditions. Developing a MATLAB code, the analytical results are validated with the results obtained using ABAQUS finite element commercial software and those available in the literature. Many parametric studies have been

manifested by applying the present analytical method and ABAQUS.

2. Material Modeling

A rectangular perforated plate lies in $(0,0,0) \leq (x, y, z) \leq (a, b, h)$ is considered as indicated in Fig. 1a. The plate consists of a repetitive pattern of rectangular voids with equal distances from each other.

Pursuant to Fig 1. it is required to define a mathematical function considering material properties in orthogonal paths. To this end, Heaviside distribution functions are introduced by Eqs. 1(a) and 1(b) [16].

$$H(x) = \sum_{i=1}^{m_x} \sum_{j=1}^{n_y} \left(\text{Heaviside}(x - x_{ci} + c/2) - \text{Heaviside}(x - x_{ci} - c/2) \right) \quad (1a)$$

$$H(y) = \sum_{i=1}^{m_x} \sum_{j=1}^{n_y} \left(\text{Heaviside}(y - y_{cj} + d/2) - \text{Heaviside}(y - y_{cj} - d/2) \right) \quad (1b)$$

where (x_{ci}, y_{ci}) is the location of the center of voids. The parameters c and d are the length & width of voids.

The Heaviside distribution (HD) function is presented by Eq. 2 [16].

$$HD = 1 - H(x) \cdot H(y) \quad (2)$$

By plotting the HD function Fig 2., it can be observed that values of one and zero are allocated to white and black regions.

The stiffness matrix of PCP layers can be given by:

$$Q(x, y)^k = Q^k \cdot HD \quad (3)$$

where Q^k is the stiffness matrix of an orthotropic lamina given by Eq. 4 [6].

$$Q^k = \begin{bmatrix} \frac{E_1}{1 - \nu_{12}\nu_{21}} & \frac{\nu_{12}E_2}{1 - \nu_{12}\nu_{21}} & 0 \\ \frac{\nu_{21}E_1}{1 - \nu_{12}\nu_{21}} & \frac{E_2}{1 - \nu_{21}\nu_{12}} & 0 \\ 0 & 0 & G_{12} \end{bmatrix} \quad (4)$$

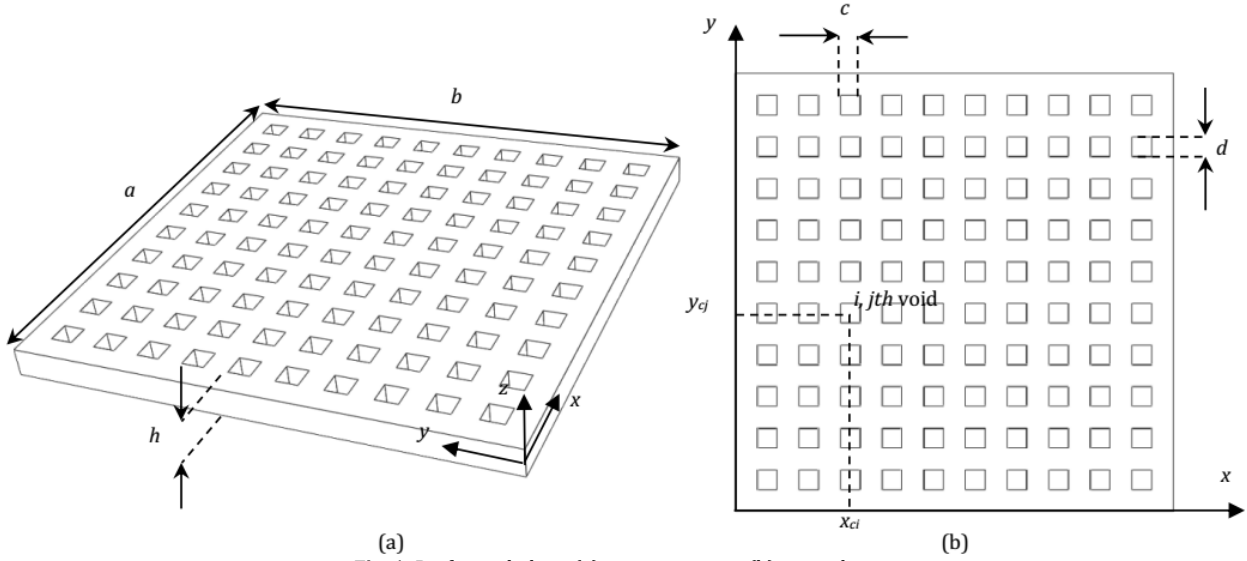


Fig. 1. Perforated plate: (a) isometric view, (b) normal view.

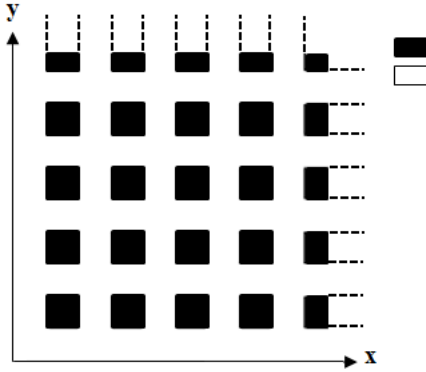


Fig. 2. The plot of the HD function.

3. Theoretical Formulations

The linear displacement field is contemplated as [6]:

$$u(x, y, z) = u_0(x, y, z) - zw_{0,x}(x, y) \quad (5)$$

$$v(x, y, z) = v_0(x, y, z) - zw_{0,y}(x, y) \quad (6)$$

$$w(x, y, z) = w_0(x, y) \quad (7)$$

Applying CLPT, equilibrium equations have been derived as [6]:

$$N_{x,x} + N_{xy,y} + N_x^i u_{0,xx} + N_y^i u_{0,yy} + N_{xy}^i u_{0,xy} = \rho h \ddot{u}_0 \quad (8)$$

$$N_{xy,x} + N_{y,y} + N_x^i v_{xx} + N_y^i v_{0,yy} + N_{xy}^i v_{0,xy} = \rho h \ddot{v}_0 \quad (9)$$

$$M_{x,xx} + 2M_{xy,xy} + M_{y,yy} + N_x^i w_{xx} + N_y^i w_{yy} + N_{xy}^i w_{xy} = \rho h \ddot{w}_0 \quad (10)$$

Force and moment resultants are given by [6]:

$$\begin{bmatrix} N_x \\ N_y \\ N_{xy} \\ M_x \\ M_y \\ M_{xy} \end{bmatrix} = \begin{bmatrix} A_{11} & A_{12} & A_{16} & B_{11} & B_{12} & B_{16} \\ A_{21} & A_{22} & A_{26} & B_{21} & B_{22} & B_{26} \\ A_{61} & A_{62} & A_{66} & B_{61} & B_{62} & B_{66} \\ B_{11} & B_{12} & B_{16} & D_{11} & D_{12} & D_{16} \\ B_{21} & B_{22} & B_{26} & D_{21} & D_{22} & D_{26} \\ B_{61} & B_{62} & B_{66} & D_{61} & D_{62} & D_{66} \end{bmatrix}$$

$$\begin{bmatrix} u_{0,x} \\ v_{0,y} \\ u_{0,y} + v_{0,x} \\ -w_{0,xx} \\ -w_{0,yy} \\ -2w_{0,xy} \end{bmatrix} \quad (11)$$

where the A , B , and D coefficients can be acquired by [6]:

$$A_{ij} = \sum_{k=1}^N (Q_{ij})_k (z_k - z_{k-1}) \quad (12)$$

$$B_{ij} = \frac{1}{2} \sum_{k=1}^N (Q_{ij})_k (z_k^2 - z_{k-1}^2) \quad (13)$$

$$D_{ij} = \frac{1}{3} \sum_{k=1}^N (Q_{ij})_k (z_k^3 - z_{k-1}^3) \quad (14)$$

And the initial uniform thermal load can be shown by [6]:

$$N_{x,y}^i = Q(x, y)_{3 \times 3}^k \alpha_{3 \times 1}^k h \Delta T \quad (15)$$

where the parameter ΔT refers to uniform temperature difference.

The governing equilibrium equations of PCPs can be obtained by substitution of Eqs. (11) and (15) into Eqs. (8-10) as:

$$\begin{aligned} & (A_{11,x} + A_{16,y})u_{0,x} + (A_{16,x} + A_{66,y})u_{0,y} \\ & + (A_{16,x} + A_{66,y})v_{0,x} + (A_{12,x} + A_{26,y})v_{0,y} \\ & - (B_{11,x} + B_{16,y})w_{0,xx} - 2(B_{16,x} + B_{66,y})w_{0,xy} \\ & - (B_{12,x} + B_{26,y})w_{0,yy} + A_{11}u_{0,xx} + 2A_{16}u_{0,xy} \\ & + A_{66}u_{0,yy} + A_{16}v_{0,xx} + (A_{12} + A_{66})v_{0,xy} + A_{26}v_{0,yy} \\ & - B_{11}w_{0,xxx} - 3B_{16}w_{0,xxxy} - (B_{12} + 2B_{66})w_{0,xxyy} \\ & - B_{26}w_{0,yyy} + N_x^i u_{xx} + N_y^i u_{yy} + N_{xy}^i u_{xy} \\ & = \rho h \ddot{u}_0 \end{aligned} \quad (16)$$

$$\begin{aligned} & (A_{16,x} + A_{12,y})u_{0,x} - (A_{66,x} + A_{26,y})u_{0,y} \\ & + (A_{66,x} + A_{26,y})v_{0,x} + (A_{26,x} + A_{22,y})v_{0,y} \\ & - (B_{16,x} + B_{12,y})w_{0,xx} - 2(B_{66,x} + B_{26,y})w_{0,xy} \\ & - (B_{26,x} + B_{22,y})w_{0,yy} + A_{16}u_{0,xx} + (A_{12} + A_{66})u_{0,xy} \\ & + A_{26}u_{0,yy} + A_{66}v_{0,xx} + 2A_{26}v_{0,xy} + A_{22}v_{0,yy} \\ & - B_{16}w_{0,xxx} - (B_{12} + 2B_{66})w_{0,xxxy} - 3B_{26}w_{0,xxyy} \\ & - B_{22}w_{0,yyy} + N_x^i v_{xx} + N_y^i v_{yy} + N_{xy}^i v_{xy} \\ & = \rho h \ddot{v}_0 \end{aligned} \quad (17)$$

$$\begin{aligned} & (B_{11,xx} + 2B_{16,xy} + B_{12,yy})u_{0,x} \\ & + (B_{16,xx} + 2B_{66,xy} + B_{26,yy})u_{0,y} \\ & + (B_{16,xx} + 2B_{66,xy} + B_{26,yy})v_{0,x} \\ & + (B_{12,xx} + 2B_{26,xy} + B_{22,yy})v_{0,y} \\ & - (D_{11,xx} + 2D_{16,xy} + D_{12,yy})w_{0,xx} \\ & - (2D_{16,xx} + 4D_{66,xy} + 2D_{26,yy})w_{0,xy} \\ & - (D_{12,xx} + 2D_{26,xy} + D_{22,yy})w_{0,yy} \\ & + 2(B_{11,x} + B_{16,y})u_{0,xx} + 2(2B_{16,x} + B_{66,y} + B_{12,y})u_{0,xy} \\ & + 2(B_{66,x} + B_{26,y})u_{0,yy} + 2(B_{16,x} + B_{66,y})v_{0,xx} \\ & + 2(B_{12,x} + 2B_{26,y} + B_{66,x})v_{0,xy} + 2(B_{26,x} + B_{22,y})v_{0,yy} \\ & - 2(D_{11,x} + D_{16,y})w_{0,xxx} \\ & - (6D_{16,x} + 2D_{12,y} + 4D_{66,y})w_{0,xxxy} \\ & - (2D_{12,x} + 6D_{26,y} + 4D_{66,x})w_{0,xxyy} \\ & - 2(D_{22,y} + D_{26,x})w_{0,yyy} + B_{11}u_{0,xxx} + 3B_{16}u_{0,xxxy} \\ & + (2B_{66} + B_{12})u_{0,xxyy} + B_{26}u_{0,yyy} + B_{16}v_{0,xxx} \\ & + (B_{12} + 2B_{66})v_{0,xxxy} + 3B_{26}v_{0,xxyy} + B_{22}v_{0,yyy} \\ & - D_{11}w_{0,xxxx} - 4D_{16}w_{0,xxxy} - (2D_{12} + 4D_{66})w_{0,xxyy} \\ & - 4D_{26}w_{0,xxyy} - D_{22}w_{0,yyyy} + N_x^i w_{xx} + N_y^i w_{yy} \\ & + N_{xy}^i w_{xy} = \rho h \ddot{w}_0 \end{aligned} \quad (18)$$

where the second indices are devoted to local derivatives.

By considering displacement functions corresponding to SSSS edge conditions as [19]:

$$\begin{aligned} u_0(x,y) & = U_{mn} \cos\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{b}y\right) T(t) \end{aligned} \quad (19)$$

$$\begin{aligned} v_0(x,y) & = V_{mn} \sin\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right) T(t) \end{aligned} \quad (20)$$

$$\begin{aligned} w_0(x,y) & = W_{mn} \sin\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{b}y\right) T(t) \end{aligned} \quad (21)$$

and applying the Galerkin method, the system of PDEs given by Eqs. (16-18) lead to Eqs. (22) and (23) for thermal buckling ($T(t)=1$) and thermal induced free vibration ($T(t)=e^{i\omega t}$) analyses, respectively. The parameters ω and t refer to frequency and time, respectively.

For thermal buckling and thermal induced free vibration analyses, the eigenvalue problems given by Eq. (22) and (23) [19] have been solved.

$$([k^e] + [k^T])\{\delta\} = 0 \quad (22)$$

$$[m]\{\ddot{\delta}\} + (k^e + k^T)\{\delta\} = 0 \quad (23)$$

where k^e and k^T refer to elastic and thermal stiffness matrix coefficients, respectively. M is the mass matrix, and δ is the displacement vector. k^e , k^T and M matrices are expanded in Appendices (1-3).

4. FEM Modeling

In order to verify the analytical results, FEM models are performed using ABAQUS to simulate thermal buckling and thermal induced free vibration problems.

The FEM meshed model Fig 3. is produced utilizing S4R elements. For a PCP taking into account the material and geometry properties according to Tables 1 and 2, mesh convergence is reported for critical temperature difference and fundamental frequency in Table 3. As indicated in Table 3, enough value for element edge length to plate edge length ratio is achieved by 0.00225.

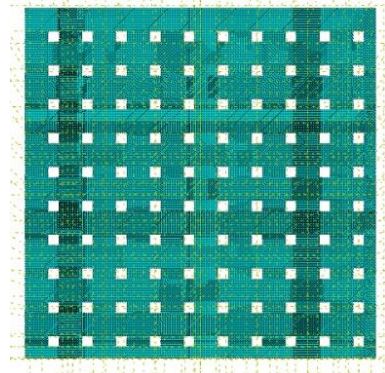


Fig. 3. FEM Meshed model.

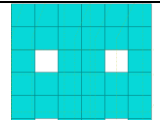

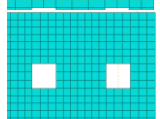
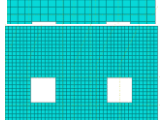
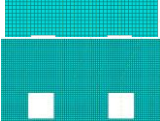
Table 1. Material properties for the PCP [20]

E_1 (GPa)	E_2 (GPa)	G_{12} (GPa)	ν_{12}	α_1 (1/°C)	α_2 (1/°C)	ρ (kg/m ³)
38.6	8.27	4.14	0.26	8.6×10^{-6}	22.1×10^{-6}	1800

Table 2. Geometry properties for the PCP

a (mm)	b (mm)	h (mm)	R
200	200	4	10%

Table 3. Mesh convergence for the PCP

element size to plate length ratio	ΔT_c (°C)	Fundamental frequency (Hz)	A partial view of the meshed model
0.02	38.228	258.62	
0.01	37.936	255.42	
0.005	38.038	255.58	
0.00225	38.072	255.18	
0.00125	38.067	255.06	

5. Validation

The accuracy of analytical solutions for buckling and free vibration are compared with those reported in the literature. For an isotropic plate ($\nu=0.3$, $\alpha=7.4 \times 10^{-6}$ 1/°C) considering different values of a/h ratio, the accuracy of the present analytical and ABAQUS methods for thermal buckling are examined through comparisons with the Ref. [14] as indicated in Table 4. A maximum discrepancy of 1.4% is acquired between the present analytical and Cetkovic results. Consequently, the present analytical method can be applied reliably for isotropic plates with a/h ratios bigger than or equal to 20.

Table 4. Comparisons of critical temperature difference results (°C) for an isotropic plate ($\nu=0.3$, $\alpha=7.4 \times 10^{-6}$ 1/°C)

Solution Method	a/h				
	20	40	60	80	100
CLPT (2002) (Adapted from [14])	427.477	106.869	47.497	26.717	17.099
Cetkovic (2016) [14]	421.584	106.497	47.423	26.694	17.089
Present Analytical	427.478	106.869	47.498	26.717	17.099
Present ABAQUS	403.99	104.27	46.828	26.435	16.964
Discrepancy between present analytical and Cetkovic results (2016)	1.4%	0.35%	0.160%	0.09%	0.06%

Considering a composite plate (layup: [0/90] s, $E_1/E_2=20$, $G_{12}/E_2=G_{13}/E_2=G_{23}/E_2=0.5$, $\nu_{12}=\nu_{13}=\nu_{23}=0.25$, $\alpha_1/\alpha_2=2$), the second validation is conducted to check the accuracy of the present analytical solution with the analytical method developed by Bouzza [13]. As indicated in Fig 4., the ABAQUS results are closer to the Ref. [13] values in comparison with the analytical results. The maximum discrepancy between the present analytical results and Ref. [13] is about 5.9%, which corresponds to a/h=2.5.

Table 5 illustrated the results of Non-dimensional frequency ($\omega^*=\omega a^2 \sqrt{\rho h / (D (2,2))}$) with corresponding mode shapes for an isotropic plate with singular central perforation ($\nu=0.3$ and $c/a=0.2$). Table 5 includes the results obtained by Kalita and Haldar [15] and the present study results and the maximum discrepancy is 3.08%. The present ABAQUS results are in closer agreement with the Ref. [15] since both methods are developed based on nonlinear displacement field.

6. Results and Discussion

In order to evaluate the thermal buckling behavior of PCPs efficiently, some case studies are demonstrated.

For examined perforated plates, the number of voids ($m_x = n_y$) are contemplated equal to 20, and the volume fraction of voids is defined as:

$$R = m_x n_y \frac{c \times d}{a \times b} \times 100 \tag{24}$$

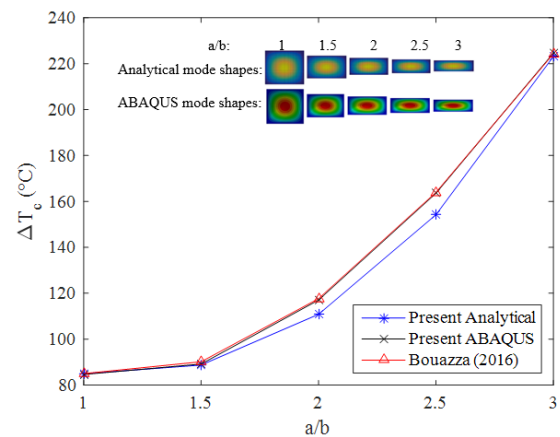
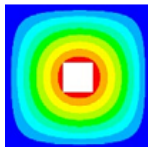
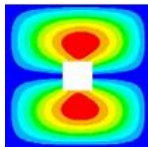
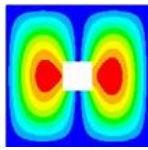
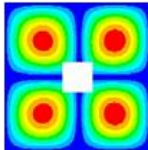
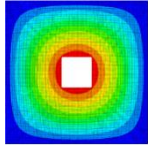
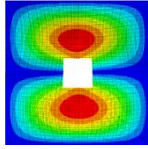
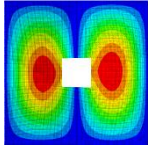
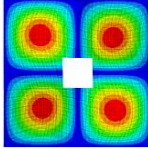
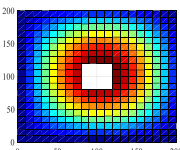
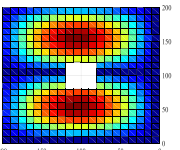
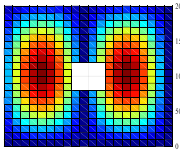
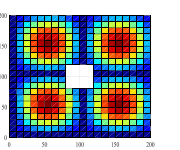


Fig. 4. Comparison of critical temperature difference results versus a/b ratio for a composite plate.

Table 5. Non-dimensional frequency ($\omega^* = \omega a^2 \sqrt{\rho h / D(2.2)}$) results for isotropic plate with singular central perforation ($\nu=0.3, c/a=1/5$)

Solution Method	Mode(1,1)	Mode(1,2)	Mode(2,1)	Mode(2,2)
	19.1	47.496	47.496	76.25
Kalita (2016) [15]				
	19.128	47.6649	47.6649	76.39
Present ABAQUS				
	18.753	48.961	48.961	77.17
Present analytical				
% Discrepancy, Present Analytical	1.82	-3.08	-3.08	-1.21
% Discrepancy, Present ABAQUS	0.15	0.36	0.36	0.18

6.1. Effect of elastic modulus and thermal expansion coefficient ratios (E_2/E_1 and α_2/α_1) on the critical temperature difference

In order to study the effect of parameters E_2/E_1 and α_2/α_1 on thermal buckling of PCPs, the material properties are shown in Table 6. The stacking sequence is assumed to be [0]. As illustrated in Fig. 5, by increasing the values of E_2/E_1 and α_2/α_1 ratios, the critical temperature difference decreases. A maximum discrepancy of 5.82% is accomplished between ABAQUS and analytical results where $E_2/E_1=0.2$ and $\alpha_2/\alpha_1=1.2$.

6.2. Comparison between the critical temperature difference of PCP and SWCCP

The critical temperature difference of PCP and the same weight monolithic composite plate (SWCCP) is indicated in Table 7. The stacking sequence is considered to be [0/90] s. As the

volume fraction of the voids increases, the PCP demonstrates higher resistance to thermal buckling than the SWCCP. It can be observed that a PCP can ameliorate the critical temperature difference four times higher than SWCCP at $R=50\%$. Figs 6(a). and 6(b). indicate the critical thermal buckling modes for PCP and SWCCP, respectively.

6.3. Effect of elastic modulus ratio (E_2/E_1) on the thermal induced fundamental frequency

In order to explore the effect of the E_2/E_1 ratio on the thermal induced fundamental frequency of PCPs, the material properties are considered as presented in Table 8.

Table 6. Material properties to study the effect of E_2/E_1 and α_2/α_1 ratios

E_1 (GPa)	E_2 (GPa)	G_{12} (GPa)	ν_{12}	α_1 (1/°C)	α_2 (1/°C)
40	variable	15.873	0.26	10^{-6}	variable

Table 7. Critical temperature difference results for PCP ($R=50\%$) and SWCCP.

R (%)	t (mm)	PCP [0/90]s		t (mm)	SWCCP [0/90]s		$\Delta T_c (SWCCP) / \Delta T_c (PCP)$	
		ΔT_c (Analytical) (°C)	ΔT_c (ABAQUS) (°C)		ΔT_c (Analytical) (°C)	ΔT_c (ABAQUS) (°C)	Analytical	ABAQUS
10	4	38.2243	38.072	3.6	30.9689	30.531	1.23	1.25
20	4	38.2153	38.216	3.2	24.4693	24.168	1.56	1.58
30	4	38.2062	37.773	2.8	18.7343	18.537	2.04	2.04
40	4	38.1971	36.773	2.4	13.764	13.642	2.78	2.7
50	4	38.188	35.279	2	9.5583	9.4896	4	3.72

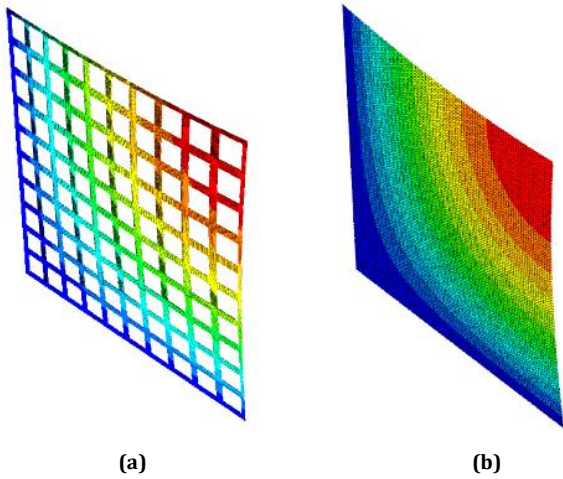


Fig. 6. Thermal buckling critical mode shape obtained by modeling a quarter of the model using ABAQUS: (a) SWCCP, (b) PCP ($R=50\%$).

The stacking sequence is assumed to be [0]. As displayed in Fig 7., by increasing the E_2/E_1 ratio, the thermal induced fundamental frequencies demonstrate ascending and descending behavior before and after the intersection point at the coordinate (32.28,272.67). A maximum discrepancy of 8.27% is achieved between ABAQUS and analytical results at $E_2/E_1=0.6$.

6.4. Effect of thermal expansion coefficient ratio (α_2/α_1) on the thermal induced fundamental frequency

In order to study the effect of the α_2/α_1 ratio on the thermal induced fundamental frequency of PCPs, the material properties are considered as indicated in Table 6. The stacking sequence is assumed [0]. As portrayed in Fig. 8, by increasing the α_2/α_1 ratio, the thermal induced fundamental frequency is decreased. A maximum discrepancy of 9.14% is observed between ABAQUS and analytical results for $T=0$ and $\alpha_2/\alpha_1=1.25$.

Table 8. Material properties to study the effect of E_2/E_1 ratio

E_1 (GPa)	E_2 (GPa)	G_{12} (GPa)	ν_{12}	α_1 ($1/^\circ\text{C}$)	α_2 ($1/^\circ\text{C}$)
40	variable	15.873	0.26	10^{-6}	10^{-6}

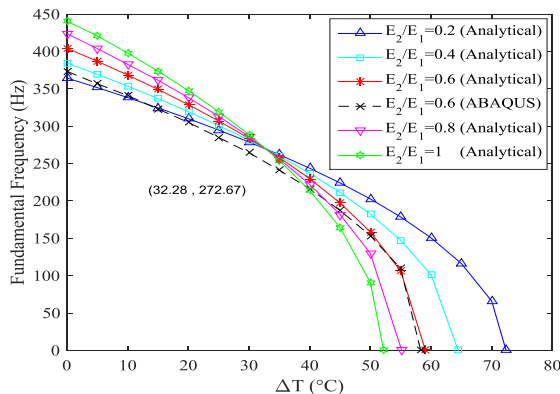


Fig. 7. Analytical results for the thermal induced fundamental frequency of PCPs considering different E_2/E_1 ratios.

6.5. Effect of stacking sequence on the thermal induced fundamental frequency

In order to study the effect of the stacking sequence on the free vibration behavior of PCPs, [0/90]_s and [0/90] stacking sequences are considered. Pursuant to Fig. 9, the gap between the thermal induced fundamental frequency corresponded to the symmetric stacking sequence, and the asymmetric one is large. It can be concluded that the PCP with symmetric stacking sequence vibrates with higher frequency. Furthermore, it can be observed that the symmetric the structure buckles at higher temperature difference than the asymmetric one.

6.6. Comparison between thermal induced fundamental frequency of PCP and SWCCP

The thermal induced fundamental frequency of PCP and SWCCP is indicated in Table 7 for the void volume fraction of $R=10\%$. The fundamental frequency of the PCP decreases with lower rate in comparison with SWCCP as the temperature difference increases

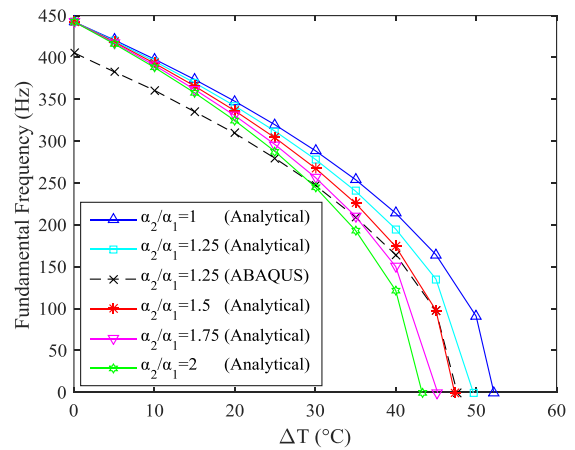


Fig. 8. Analytical results for the thermal induced fundamental frequency of PCPs considering different α_2/α_1 ratios.

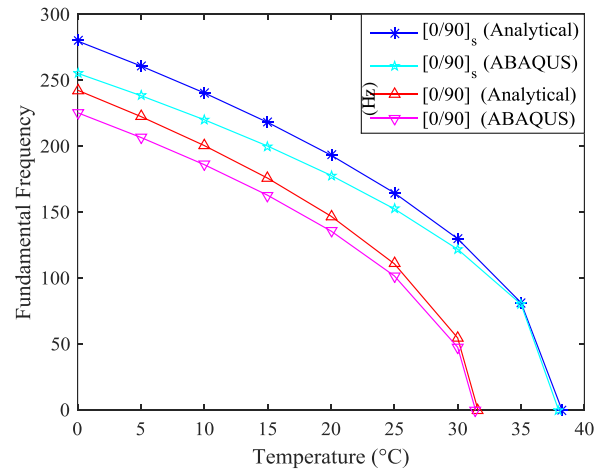


Fig. 9. Analytical results for the thermal induced fundamental frequency of PCPs considering different α_2/α_1 ratios.

7. Conclusions

A mathematical approach to thermal buckling and thermal induced free vibration analyses of perforated composite plates (PCPs) are described and discussed in this research. The present results are compared to the previously published researches [13-15], and they were found to be in close agreement.

Pursuant to results, several concluding remarks are listed as follows:

- By increasing the E2/E1 and α_2/α_1 ratios, the critical temperature difference of the orthotropic perforated plate is reduced.
- Applying the same material properties, the PCP (R=50%) is 4 times more resistant for thermal buckling than the monolithic composite plate of the same weight.
- By increasing the E2/E1 ratio, the thermal induced fundamental frequency of the orthotropic perforated plate operates ascending and descending behaviour before and after a specific intersection point.
- It can be concluded that the PCP with symmetric stacking sequence vibrates with higher frequency and buckles at higher critical temperature difference rather than the PCP with an asymmetric layup.
- Using the same material properties, variation of the fundamental frequency of the PCP (at R=10%) is less than the monolithic composite plate of the same weight by boosting the temperature.

Appendix 1

$$k_{11}^e = \int_0^b \int_0^a ((A_{11.x} + A_{16.y})u_{0,x} + (A_{16.x} + A_{66.y})u_{0,y} + A_{16}u_{0,xx} + (A_{12} + A_{66})u_{0,xy} + A_{26}u_{0,yy})U_{mn}(x,y)dxdy \quad (A.1)$$

$$k_{12}^e = \int_0^b \int_0^a ((A_{16.x} + A_{66.y})v_{0,x} + (A_{12.x} + A_{26.y})v_{0,y} + A_{16}v_{0,xx} + (A_{12} + A_{66})v_{0,xy} + A_{26}v_{0,yy})V_{mn}(x,y)dxdy \quad (A.2)$$

$$k_{13}^e = \int_0^b \int_0^a (-(B_{11.x} + B_{16.y})w_{0,xx} - 2(B_{16.x} + B_{66.y})w_{0,xy} - (B_{12.x} + B_{26.y})w_{0,yy} - B_{11}w_{0,xxx} - 3B_{16}w_{0,xyy} - (B_{12} + 2B_{66})w_{0,xyy} - B_{26}w_{0,yyy})W_{mn}(x,y)dxdy \quad (A.3)$$

$$k_{21}^e = \int_0^b \int_0^a ((A_{16.x} + A_{12.y})u_{0,x} - (A_{66.x} + A_{26.y})u_{0,y} + A_{16}u_{0,xx} + (A_{12} + A_{66})u_{0,xy} + A_{26}u_{0,yy})U_{mn}(x,y)dxdy \quad (A.4)$$

$$k_{22}^e = \int_0^b \int_0^a ((A_{66.x} + A_{26.y})v_{0,x} + (A_{26.x} + A_{22.y})v_{0,y} + A_{66}v_{0,xx} + 2A_{26}v_{0,xy} + A_{22}v_{0,yy})V_{mn}(x,y)dxdy \quad (A.5)$$

$$k_{23}^e = \int_0^b \int_0^a (-(B_{16.x} + B_{12.y})w_{0,xx} - 2(B_{66.x} + B_{26.y})w_{0,xy} - (B_{26.x} + B_{22.y})w_{0,yy} - B_{16}w_{0,xxx} - (B_{12} + 2B_{66})w_{0,xyy} - 3B_{26}w_{0,xyy} - B_{22}w_{0,yyy})W_{mn}(x,y)dxdy \quad (A.6)$$

$$k_{31}^e = \int_0^b \int_0^a ((B_{11.xx} + 2B_{16.xy} + B_{12.yy})u_{0,x} + (B_{16.xx} + 2B_{66.xy} + B_{26.yy})u_{0,y} + 2(B_{11.x} + B_{16.y})u_{0,xx} + 2(2B_{16.x} + B_{66.y} + B_{12.y})u_{0,xy} + 2(B_{66.x} + B_{26.y})u_{0,yy} + B_{11}u_{0,xxx} + 3B_{16}u_{0,xyy} + (2B_{66} + B_{12})u_{0,xyy} + B_{26}u_{0,yyy})U_{mn}(x,y)dxdy \quad (A.7)$$

$$k_{32}^e = \int_0^b \int_0^a ((B_{16.x.x} + 2B_{66.x.y} + B_{26.y.y})v_{0,x} + (B_{12.xx} + B_{26.xy} + B_{22.yy})v_{0,y} + 2(B_{16.x} + B_{66.y})v_{0,xx} + 2(B_{12.xx} + 2B_{26.y} + B_{66.x})v_{0,xy} + 2(B_{26.x} + B_{22.y})v_{0,yy} + B_{16}v_{0,xxx} + (B_{12} + 2B_{66})v_{0,xyy} + 3B_{26}v_{0,xyy} + B_{22}v_{0,yyy})V_{mn}(x,y)dxdy \quad (A.8)$$

$$k_{33}^e = \int_0^b \int_0^a (-(D_{11.xx} + 2D_{16.xy} + D_{12.yy})w_{0,xx} - (2D_{16.xx} + 4D_{66.xy} + 2D_{26.yy})w_{0,xy} - (D_{12.xx} + 2D_{26.xy} + D_{22.yy})w_{0,yy} - 2(D_{11.x} + D_{16.y})w_{0,xxx} - 2(D_{11.x} + D_{16.y})w_{0,xxx} - (6D_{16.x} + 2D_{12.y} + 4D_{66.y})w_{0,xyy} - (2D_{12.x} + 6D_{26.y} + 4D_{66.x})w_{0,xyy} - 2(D_{22.y} + D_{26.x})w_{0,yyy} - D_{11}w_{0,xxxx} - (2D_{12} + 4D_{66})w_{0,xyyy} - 4D_{26}w_{0,xyyy} - D_{22}w_{0,yyyy})W_{mn}(x,y)dxdy \quad (A.9)$$

Appendix 2

$$k_{11}^T = \int_0^a \int_0^b (N_x^i u_{xx} + N_y^i u_{yy})U_{mn}(xy)dxdy$$

$$k_{12}^T = \int_0^a \int_0^b N_{xy}^i u_{xy}V_{mn}(xy)dxdy$$

$$k_{21}^T = \int_0^a \int_0^b N_{yx}^i u_{yx} U_{mn}(xy) dx dy$$

$$k_{22}^T = \int_0^a \int_0^b (N_x^i v_{xx} + N_y^i v_{yy}) V_{mn}(xy) dx dy$$

$$k_{23}^T = 0$$

$$k_{31}^T = 0$$

$$k_{32}^T = 0$$

$$k_{33}^T = \int_0^a \int_0^b (N_x^i w_{xx} + N_y^i w_{yy}) W_{mn}(xy) dx dy$$

Appendix 3

$$M_{11} = \int_0^a \int_0^b \int_{-h/2}^{h/2} \rho U_{mn}(xy) dx dy$$

$$M_{12} = 0$$

$$M_{13} = 0$$

$$M_{21} = 0$$

$$M_{22} = \int_0^a \int_0^b \int_{-h/2}^{h/2} \rho V_{mn}(xy) dx dy$$

$$M_{31} = 0$$

$$M_{32} = 0$$

$$M_{33} = \int_0^a \int_0^b \int_{-h/2}^{h/2} \rho W_{mn}(xy) dx dy$$

References

[1] Gossard ML, Seide P, Roberts WM. Thermal Buckling of Plates, NACA TND. USA; 1995.

[2] Tauchert TR, Huang N N. Thermal Buckling of Symmetric Angle-ply Laminated Plates. *Compos. Struct* 1987; 4: 424-35.

[3] Thangaratnam KR, Palaninathan, Ramachandaran J. Thermal Buckling of Composite Laminated Plates. *Comput. Struct* 1989; 32: 1117-24.

[4] Sunt LX, Hsij TR. Thermal Buckling of Laminated Composite Plates with Transverse Shear Deformation. *Comput. Struct* 1990; 36: 883-89.

[5] Thornton EA. Thermal Buckling of Plates and Shells. *Appl. Mech* 1993; 46: 485-06.

[6] Whitney JM. Structural Analysis of Laminated Anisotropic Plates. Lancaster, PA: Technomic Publishing Co. Inc; 1987.

[7] Kant T, Babu CS. Thermal Buckling Analysis of Skew Fiber-Reinforced Composite and Sandwich Plates Using Shear Deformable Finite Element Models. *Compos. Struct* 2000; 49: 77-85.

[8] Yapici A. Thermal Buckling Behavior of Hybrid-Composite Angle-Ply Laminated Plates with an Inclined Crack. *Mech. Compos. Mat* 2005; 41: 131-38.

[9] Pradeep V, Ganesan N. Thermal Buckling and Vibration Behavior of Multi-Layer Rectangular Viscoelastic Sandwich Plates. *J. Sound. Vib* 2008; 310: 169 – 83.

[10] Duran AV, Fasanella NA, Sundararaghavan V, Waas A. Thermal Buckling of Composite Plates with Spatial Varying Fiber Orientations, *Compos. Struct* 2015; 124: 228-35.

[11] Li X, Yu K, Han J, Song H, Zhao R. Buckling and Vibro-Acoustic Response of the Clamped Composite Laminated Plate in Thermal Environment. *Int. J. Mech. Sci* 2016; 119: 370-82.

[12] Jin T, Ha NS, Le VT, Goo NS, Thermal Buckling Measurement of a Laminated Composite Plate Under a Uniform Temperature Distribution Using the Digital Image Correlation Method. *Compos. Struct* 2015; 123: 420-2.

[13] Bouzza M, Lairedj A, Benseddiq N, Khalki S. a Refined Hyperbolic Shear Deformation Theory for Thermal Buckling Analysis of Cross-Ply Laminated Plates. *Mech. Res. Commun* 2016; 73: 117-26.

[14] Cetkovic M, Thermal Buckling of Laminated Composite Plates Using Layerwise Displacement Model, *Compos. Struct* 2016; 142: 238-53.

[15] Kalita K, Halder S. Free Vibration Analysis of Rectangular Plates with Central Cutout. *Cogent Eng* 2016; 3: 1163781.

[16] Takabatake H. Static Analyses of Elastic Plates with Voids. *Int. J. Solids. Struct* 1991; 28: 179-96.

[17] Li G, Cheng J. A Generalized Analytical Modeling of Grid Stiffened Composite Structures. *Compos. Struct* 2012; 94: 1117–27.

[18] Wilson AJ, Rajasekaran S. Stability of All Edges Simply Supported, Stepped and Stiffened Rectangular Plate Under Biaxial Loading. *Appl. Math. Modeling* 2014; 38: 479–95.

[19] Jones R M. Buckling of Bars, Plates and Shells. Bull Ridge Corporation; 2006.

[20] Kaw A K. Mechanics of Composite Materials, Second Edition. New York: CRC Press, Taylor & Francis Group 2006.